NE150/215M Introduction to Nuclear Reactor Theory Spring 2022

Discussion 5: Neutron Transport

February 23, 2022

Helpful Readings: LE Ch. 6, LB Ch. 5.

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The rate of a reaction can be calculated as:

$$R = V\Phi\Sigma_{\text{reaction}}$$
 $\left[\frac{\text{reaction}}{\text{second}}\right]$

- 1. Use this to derive the heterogenous expression for the:
 - a. Thermal neutron utilization factor
 - b. Thermal neutron reproduction factor



- 1. Derive the heterogenous expression for the:
 - a. Thermal neutron utilization factor

$$f = \frac{Thermal\ neutrons\ absorbed\ in\ fuel}{Thermal\ neutrons\ absorbed\ in\ all\ reactor\ materials}$$



- 1. Derive the heterogenous expression for the:
 - a. Thermal neutron utilization factor

$$f = \frac{Thermal\ neutrons\ absorbed\ in\ fuel}{Thermal\ neutrons\ absorbed\ in\ all\ reactor\ materials} \\ f = \frac{R_{abs}^f}{R_{abs}^f + R_{abs}^m + R_{abs}^{oth}} = \frac{V^f \overline{\Phi}^f \overline{\Sigma}_{abs}^f}{V^f \overline{\Phi}^f \overline{\Sigma}_{abs}^f + V^m \overline{\Phi}^m \overline{\Sigma}_{abs}^m + V^{oth} \overline{\Phi}^{oth} \overline{\Sigma}_{abs}^{oth}}$$

Note, flux and macroscopic cross sections are averaged over space and energy.



- 1. Derive the heterogenous expression for the:
 - b. Thermal neutron reproduction factor

$$\eta = \frac{\textit{Neutrons produced by fission}}{\textit{Thermal neutrons absorbed in the fuel}}$$



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 - b. Thermal neutron reproduction factor

$$\eta = \frac{\textit{Neutrons produced by fission}}{\textit{Thermal neutrons absorbed in the fuel}}$$

$$\eta(E) = \frac{\nu R_{fis}^f}{R_{abs}^f} = \nu \frac{V^f \overline{\Phi}^f \overline{\Sigma}_{fis}^f}{V^f \overline{\Phi}^f \overline{\Sigma}_{abs}^f}$$

Note, flux and macroscopic cross sections are averaged over space and energy.



Review



Calculating k_{eff} for real reactors

For heterogenous reactors, some of our terms depend on average neutron flux and volume. How do we calculate flux?

$$f = \frac{\mathbf{V}^{f} \overline{\Phi}^{f} \overline{\Sigma}_{abs}^{f}}{\mathbf{V}^{f} \overline{\Phi}^{f} \overline{\Sigma}_{abs}^{f} + \mathbf{V}^{m} \overline{\Phi}^{m} \overline{\Sigma}_{abs}^{m} + \mathbf{V}^{oth} \overline{\Phi}^{oth} \overline{\Sigma}_{abs}^{oth}}$$

$$\eta = \nu \frac{V^{f} \overline{\Phi}^{f} \overline{\Sigma}_{fis}^{f}}{V^{f} \overline{\Phi}^{f} \overline{\Sigma}_{abs}^{f}}$$



How are neutrons distributed?

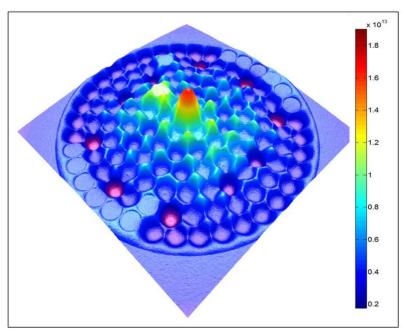
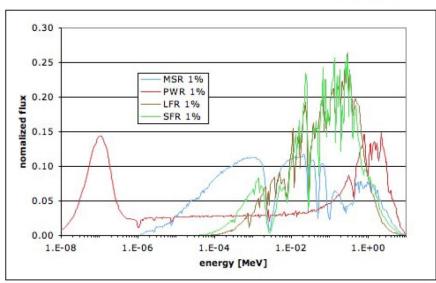


Figure 13. MCNP simulation of thermal neutron flux radial distribution at 1 MW core power level.

Mohamad Hairie B. Rabir et al 2018 IOP Conf. Ser.: Mater. Sci.

Eng. 298 012029



"Performance of molten salt versus solid fuel reactors" B. Becker, M. Fratoni, E. Greenspan



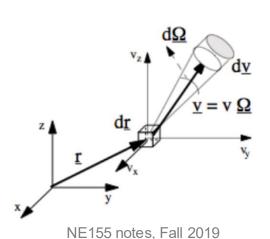
Neutron Transport Equation

- Describes neutron population in the core
- Neutron balance equation that conserves neutrons
- Each term considers either the gain or loss of neutrons



Neutron Transport Equation

$$\underbrace{\frac{1}{v}\frac{\partial\psi}{\partial t}(\vec{r},E,\hat{\Omega},t)}_{\text{time rate of change}} + \underbrace{\hat{\Omega}\cdot\nabla\psi(\vec{r},E,\hat{\Omega},t)}_{\text{streaming loss rate}} + \underbrace{\Sigma_t(\vec{r},E)\psi(\vec{r},E,\hat{\Omega},t)}_{\text{total interaction loss rate}}$$



$$=\underbrace{\int_{0}^{\infty}\int_{4\pi}\Sigma_{s}(\vec{r},E'\to E,\hat{\Omega}'\to\hat{\Omega})\psi(\vec{r},E',\hat{\Omega}',t)d\hat{\Omega}'dE'}_{\text{in scattering source rate}} \\ +\underbrace{\frac{\chi_{p}(E)}{4\pi}\int_{0}^{\infty}\int_{4\pi}\nu(E')\Sigma_{f}(\vec{r},E')\psi(\vec{r},E',\hat{\Omega}',t)d\hat{\Omega}'dE'}_{\text{fission source rate}}$$

external source rate

- 1. Particles can be treated essentially as points
 - Their state is described fully by their location, velocity vector, and a given time.
 - Rotation and quantum effects are ignored.



- Particles travel in straight lines between collisions
 - Since neutrons are neutral, you do not have to account for other forces



- 3. Collisions between particles are negligible
 - The transport equation becomes linear



- 4. Material properties are isotropic
 - Usually valid unless velocities are very low



- 5. Material properties are time-independent
 - Valid for short time scales (i.e. no depletion)



- 6. Quantities are expected values
 - Fluctuations about the mean for very low densities are not accounted for.



Solving the transport equation

- Using deterministic methods
 - Fast, but complicated
- Monte-Carlo simulations (MCNP, Serpent)
 - Expensive to run for statistically accurate results (many particles needed)
 - Stochastic
- Using the diffusion equation as an approximation
 - Simplifies the problem so it can be solved by hand
 - Not always applicable



Angular Flux

The path length per unit volume about \vec{r} passed by neutrons with energies in dE about E at time t.

$$\psi(\vec{r}, E, \widehat{\Omega}, t) \equiv vN(\vec{r}, E, \widehat{\Omega}, t)$$

units:
$$\frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV} \cdot \text{steradian}}$$



Scalar Flux

The number of neutrons penetrating a sphere with a cross sectional area of 1 cm² at \vec{r} , with energies in dE about E at time t.

$$\phi(\vec{r}, E, t) \equiv vN(\vec{r}, E, t)$$

$$\phi(\vec{r}, E, t) = \int_{4\pi} \psi(\vec{r}, E, \widehat{\Omega}, t) d\widehat{\Omega}$$

units:
$$\frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV}}$$



Net Current

Net number of particles crossing a unit area per second along a **direction normal** to that area with energies in [E, E+dE] at time t.

$$\vec{J}(\vec{r}, E, t) = \int_{4\pi} \hat{k} \cdot \widehat{\Omega} \psi(\vec{r}, E, \widehat{\Omega}, t) d\widehat{\Omega}$$

units:
$$\frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV}}$$



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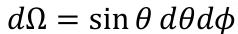
In just the positive direction, over all energies:

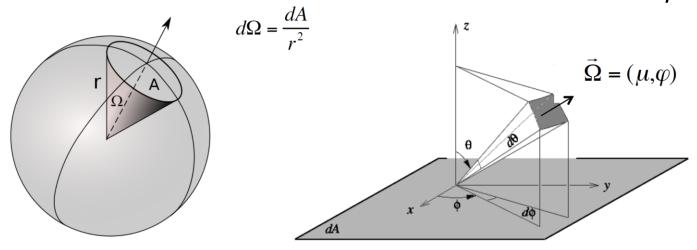
$$J^{+}(z) = \int_{2\pi^{+}} \hat{k} \cdot \widehat{\Omega} \psi(\vec{r}, \widehat{\Omega}) d\widehat{\Omega}$$

units:
$$\frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV}}$$



Solid Angle Review





"NE:150 Spring 2018 Lecture 19: Neutron Diffusion Equation – 3", J. Vujic



Practice



Recall that a neutron's energy after scattering is given by

$$E' = \frac{E}{(1+A)^2} \left(\mu + \sqrt{A^2 + \mu^2 - 1}\right)^2$$
, where $\mu = \cos \theta$

1. Write an equation for the expected value of μ (or $\bar{\mu}$)

Hint: In statistics,
$$E(X) = \bar{X} = \int_{-\infty}^{\infty} XP(X) dx$$

2. How can this be related to E' as given above?



1. Write an equation for the expected value of $\mu = \cos \theta$ Hint: In statistics, $\overline{X} = \int XP(X)dx$.



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What are the bounds of the integral?



1. Write an equation for the expected value of $\mu = \cos \theta$

Hint: In statistics, $\overline{X} = \int XP(X)dx$.

$$\bar{\mu} = \int \mu P(\mu) d\mu$$

What are the bounds of the integral?

$$\bar{\mu} = \int_{-1}^{1} \mu P(\mu) d\mu$$



$$E' = \frac{E}{(1+A)^2} \Big(\mu + \sqrt{A^2 + \mu^2 - 1}\Big)^2 \text{ , where } \mu = \cos\theta$$

$$\bar{\mu} = \int_{-1}^1 \mu P(\mu) d\mu$$

2. How can this be related to E' as given above?



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$$\bar{\mu} = \int_{-1}^1 \mu P(\mu) d\mu$$

2. How can this be related to E' as given above?

$$P(\mu)d\mu = P(E')dE'$$



In the transport equation, we neglect:

- collisions between neutrons
- 2. the impact of gravity on neutrons Show that these effects are negligible in a thermal reactor, given a scalar flux of $10^{12} \frac{n}{cm^2s}$ and water $(\rho = 1 \frac{g}{cm^3}, \ \Sigma_a = 0.0222 \ cm^{-1})$ as a moderator.



1. Show that collisions between neutrons are negligible in a thermal reactor, given a scalar flux of $10^{12} \frac{n}{cm^2s}$ and water

For thermal neutrons, $v = 2200 \frac{m}{s}$



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For thermal neutrons, $v = 2200 \frac{m}{s}$

$$N = \frac{\phi}{v} = \frac{10^{12} \ n/cm^2 s}{2200 \ m/s} \approx 4.5 \times 10^6 \frac{neutrons}{cm^3}$$



1. Show that collisions between neutrons are negligible in a thermal reactor, given a scalar flux of $10^{12} \frac{n}{cm^2s}$ and water

$$N_{neutrons} \approx 4.5 \times 10^6 \frac{neutrons}{cm^3}$$

Consider the number of density of water

$$N_{water} = \frac{\left(1\frac{g}{cm^3}\right)6.022 \times 10^{23} \frac{atom}{mol}}{18 \ gmol^{-1}} \approx 3.3 \times 10^{22} \frac{molecules}{cm^3}$$



1. Show that collisions between neutrons are negligible in a thermal reactor, given a scalar flux of $10^{12} \frac{n}{cm^2s}$ and water

$$N_{neutrons} \approx 4.5 \times 10^6 \frac{neutrons}{cm^3}, \qquad N_{water} \approx 3.3 \times 10^{22} \frac{molecules}{cm^3}$$
 $N_{neutrons} \ll N_{water}$

We don't know the cross section for neutron to neutron interactions. But, with how much bigger N_{water} is, the probability of a neutron interacting with water is FAR bigger.



2. Show that the impact of gravity on neutrons is negligible in a thermal reactor, given water as a moderator ($\Sigma_a = 0.0222 \ cm^{-1}$)



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time between absorptions =
$$t = \frac{1}{\sum_{a}^{water} v}$$

$$t = \frac{1}{(0.0222cm^{-1})(2200\frac{m}{s})} = 2.05 \times 10^{-4} s$$



2. Show that the impact of gravity on neutrons is negligible in a thermal reactor, given water as a moderator ($\Sigma_a = 0.0222 \ cm^{-1}$)

$$t = 2.05 \times 10^{-4} \text{ s}$$

displacement due to gravity = $\Delta x = \frac{1}{2}gt^2$
 $\Delta x = \frac{1}{2}(9.8 \frac{\text{m}}{\text{s}^2})(2.05 \times 10^{-4})^2 = 2.059 \times 10^{-7} \text{ m}$



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Much, much smaller than the mean free path of neutrons in water!





$$\Omega = \int d\Omega = \int_0^{\pi/4} \int_0^{\pi} \sin\theta \, d\theta d\phi$$



$$\Omega = \int d\Omega = \int_0^{\pi/4} \int_0^{\pi} \sin \theta \, d\theta d\phi$$
$$= \int_0^{\pi/4} [-\cos \theta]_0^{\pi} d\phi = \int_0^{\pi/4} [1+1] d\phi$$



$$\Omega = \int d\Omega = \int_0^{\pi/4} \int_0^{\pi} \sin\theta \, d\theta d\phi$$

$$= \int_0^{\pi/4} [-\cos\theta]_0^{\pi} d\phi = \int_0^{\pi/4} [1+1] d\phi$$

$$\Omega = 2 \int_0^{\pi/4} d\phi = \frac{\pi}{2} steradians$$

