

NE150/215M Introduction to Nuclear Reactor Theory
Spring 2022

Midterm 1 Review

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Ian Kolaja



Nuclear Physics and Reactions

Number Densities

For isotope with abundance γ_i

$$N\left({}_Z^AX\right) = \frac{\gamma_i \rho\left({}_Z^AX\right)}{M\left({}_Z^AX\right)} N_A, \quad \left(\frac{\# \text{ atoms}}{\text{cm}^3}\right)$$

For chemical X_mY_n

$$N_X = mN_{X_mY_n} \qquad N_Y = nN_{X_mY_n}$$

$$\text{Weight fraction: } w_X = \frac{mM_X}{mM_X + nM_Y}$$

Number Densities

- Review first and second homework
- Know how to handle number densities for molecules, different enrichment

$$N_{total} = \sum_i N_i = \sum_i \frac{w_i \rho N_A}{M_i}$$

$$\text{Atomic mass of mixture: } \frac{1}{M} = \sum_i \frac{w_i}{M_i}$$

Activity

Decay constants and half lives:

$$t_{1/2} = \frac{\ln 2}{\lambda}$$

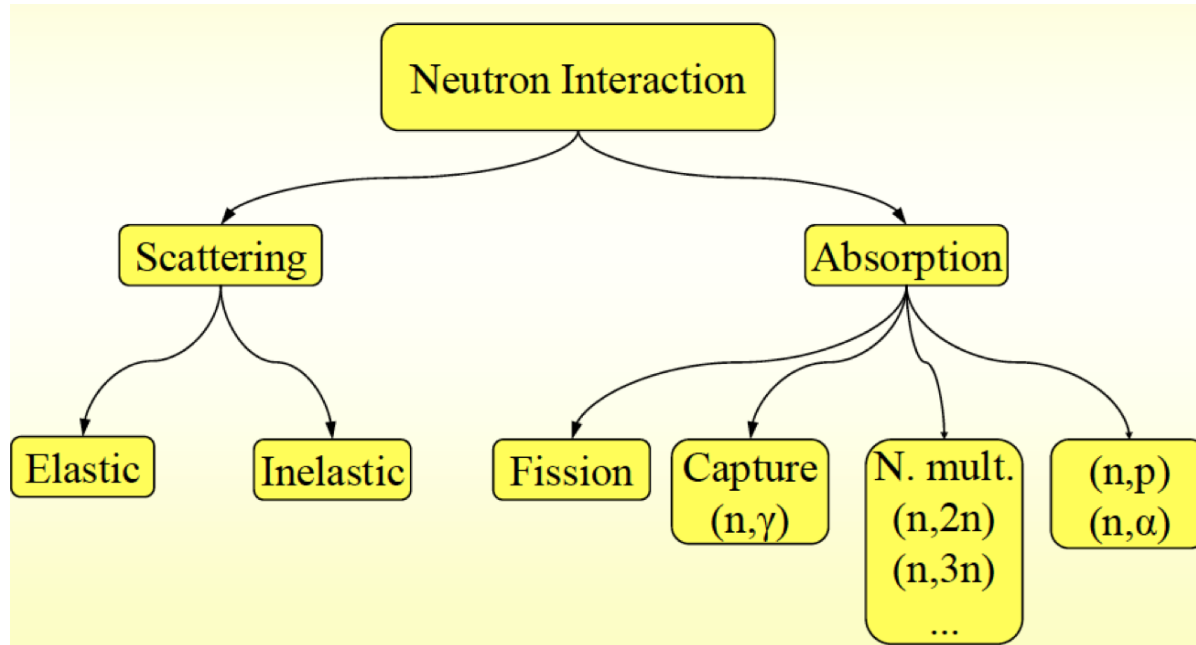
Relationship between activity and population:

$$A(t) = \lambda N(t)$$

Population over time:

$$N(t) = N_0 e^{-\lambda t}$$

Types of Reactions



“NE:150 Spring 2018 Lecture 5: Neutron Interactions with Matter”, J. Vujic

Types of Reactions

Reaction Name	Notation	Reaction	
Elastic Scattering	(n, n)	${}_0^1n + {}_Z^AX \rightarrow {}_0^1n + {}_Z^AX$	Scattering
Inelastic Scattering	(n, n')	${}_0^1n + {}_Z^AX \rightarrow ({}_Z^{A+1}X^*) \rightarrow {}_0^1n + {}_Z^AX + \gamma$	
Radiative Capture	(n, γ)	${}_0^1n + {}_Z^AX \rightarrow ({}_Z^{A+1}X^*) \rightarrow {}_Z^{A+1}X + \gamma$	
Fission	(n, f)	${}_0^1n + {}_Z^AX \rightarrow ({}_Z^{A+1}X^*) \rightarrow {}_{Z_1}^{A_1}Y + {}_{Z_2}^{A_2}Z + \nu_0^1n + \gamma$	Absorption
Charged particle reactions	(n, α)	${}_0^1n + {}_Z^AX \rightarrow {}_{Z-2}^{A-3}Y + \alpha$	
	(n, p)	${}_0^1n + {}_Z^AX \rightarrow {}_{Z-1}^AY + {}_1^1p$	
Neutron-producing reaction	$(n, 2n)$	${}_0^1n + {}_Z^AX \rightarrow {}_Z^{A-1}Y + 2{}_0^1n$	

Cross Sections - Definitions

Microscopic Cross Section (σ) – A piece of measured data that is related to the probability that a neutron will interact with a target nucleus. It's expressed as an effective area of a target, and typically measured in barns ($1 \text{ barn} = 10^{-24} \text{ cm}^2$)

Macroscopic Cross Section (Σ) – Accounts for the number density of the target nuclei, given by $\Sigma = N\sigma \text{ (cm}^{-1}\text{)}$. Conceptually, it tells you the probability per unit length that a neutron will have the given interaction.

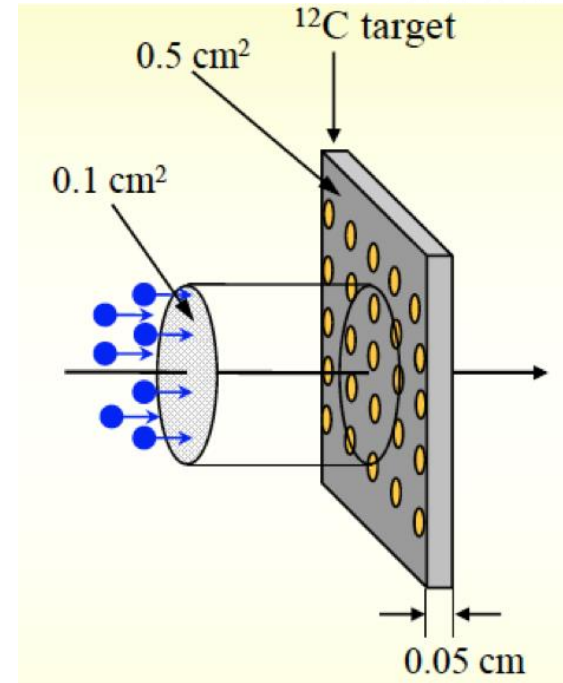
Interactions and Attenuation

- Mean free path ($\bar{x}$) – The average distance a neutron travels between collisions, given by: $\bar{x} = \frac{1}{\Sigma}$
- First collision probability: $p(x)dx = \Sigma_t e^{-\Sigma_t x} dx$
- Probability not interacting within distance x : $P(x) = e^{-\Sigma_t x}$
- Probability of at least one collision in x : $P(x)^C = 1 - e^{-\Sigma_t x}$
- Intensity of beam of uncollided neutrons after penetrating a distance of x into a material with total cross section:

$$I(x) = I_0 e^{-\Sigma_t x}$$

Reaction Rates

- Reaction rates can be calculated by determining the total number of collisions between incoming particles and nuclei per unit time and unit volume.
- Multiplying by volume gives you total reaction rate.
- Dimensional analysis is key



“NE:150 Spring 2018 Lecture 5: Neutron Interactions with Matter”, J. Vujic

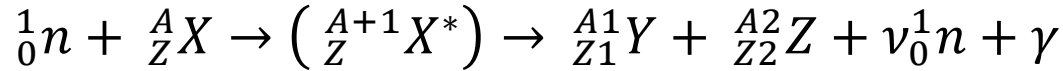
Reaction Rates

$$R = nvN\sigma = \Phi\Sigma \text{ (# Collisions/cm}^3\text{)}$$

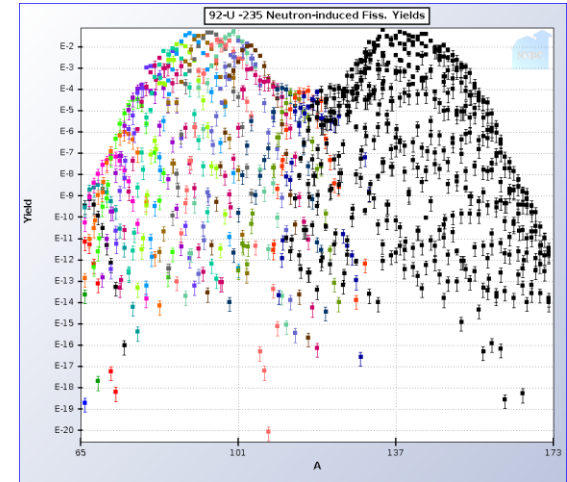
Quantity	Definition	Units
N	Target nuclei number density	<i>nucleus/cm</i> ³
n	Neutron number density	<i>neutron/cm</i> ³
v	Neutron speed	<i>m/s</i>
σ	Microscopic Cross Section	<i>cm</i> ² / <i>nucleus</i>
$\Sigma = N\sigma$	Macroscopic Cross Section	<i>1/cm</i>
$\phi = vn$	Scalar Neutron Flux	<i># neutrons/cm</i> ² <i>s</i>

Fission and Criticality

Fission



- Typically produces around 200 MeV in total
 - Depends on isotope, includes energy from decays of fission products
- Fissile material: Readily undergoes fission with any neutron energy
- Fissionable: Can undergo fission
- Fertile material: Can become fissile via neutron capture and decay
- ν , neutron produced from fission, depends on energy. Typically 2.5 for thermal fission U-235



Neutron Multiplication Factor

$$k = \frac{\text{Production rate of neutrons from fission}}{\text{loss rate of neutrons from leakage and absorption}}$$

- Subcriticality ($k < 1$)
 - Neutron population & power decrease
- Criticality ($k = 1$)
 - Chain reaction is time independent
 - Desired for reactor operation
- Supercriticality ($k > 1$)
 - Neutron population & power increase

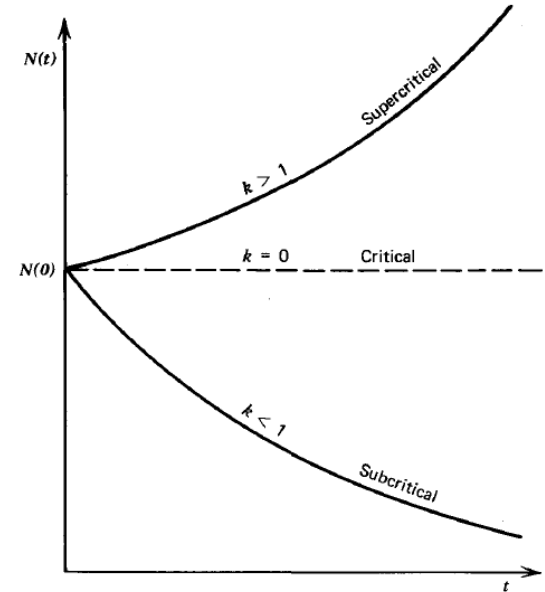


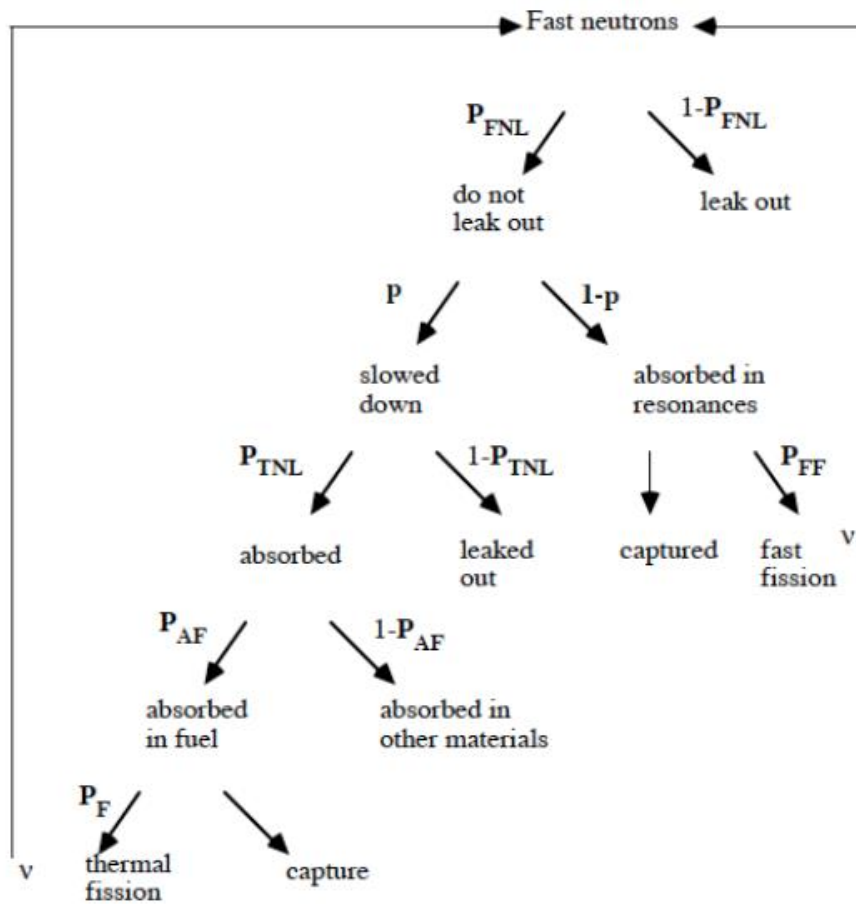
FIGURE 3-2. Time behavior of the number of neutrons in a reactor.

Duderstadt, Hamilton 1976

Six Factor Formula Terms

$$k_{eff} = \epsilon p f \eta P_{FNL} P_{TNL}$$

Var.	Name	Definition
ϵ	Fast fission factor	The total number of neutrons produced (from both thermal and fast fissions) divided by just the number of thermal fissions.
p	Resonance escape probability	The probability that a neutron passes through the resonance region (see cross section plot) without being absorbed (<1)
f	Thermal neutron utilization factor	The fraction of thermal neutrons that are absorbed in fuel materials over those absorbed in all materials.
η	Thermal neutron reproduction factor	The ratio of neutrons produced by fission over the number of thermal neutrons absorbed in the fuel.
P_{FNL}	Fast neutron non-leakage probability	The probability that a fast neutron does not leak out of the reactor
P_{TNL}	Thermal neutron non-leakage probability	The probability that a thermal neutron does not leak out of the reactor



Fuel Conversion/Breeding

- Fertile materials in the reactor can capture neutrons and decay into fissile materials.
- This can prolong your reactor operation before refueling
- Breeder reactor designs exploit this to create more fissionable material than they consume

$$C = \frac{\text{Production rate of fissile material}}{\text{Consumption rate of fissile material}} \rightarrow \frac{\text{Absorption in } U^{238}}{\text{Absorption in } U^{235}}$$

Called “Breeding Ratio, B if greater than 1.



Scattering and Neutron Transport

Elastic Scattering

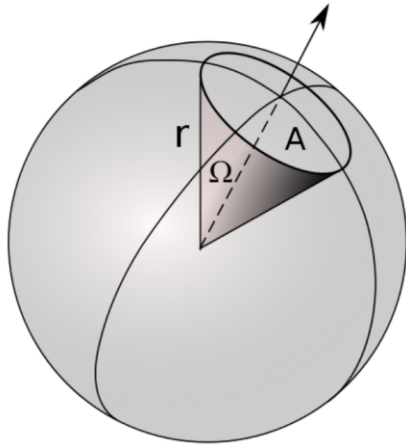
- We want to slow down neutrons because U-235 has a higher fission cross section at low energies
- Elastic scattering is the main process for slowing neutrons
 $(E, \Omega) \rightarrow (E', \Omega')$

$$\sigma_s(E \rightarrow E') = \frac{\sigma_s(E)}{(1 - \alpha)E}, \quad \alpha E \leq E' \leq E$$
$$\bar{E}' = \left(\frac{1 + \alpha}{2} \right) E, \quad \alpha = \left(\frac{A - 1}{A + 1} \right)^2$$

Elastic Scattering Properties

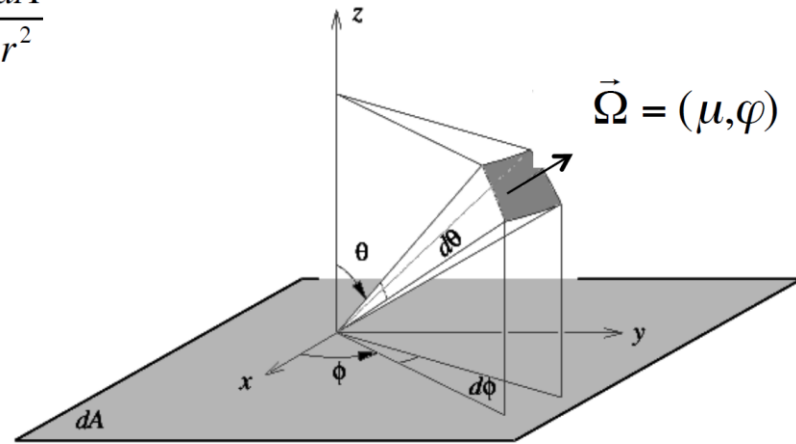
- Elastic scattering cross sections for light nuclei are nearly independent of neutron energy up to 1 MeV
- Typical cross sections for scattering range for 2 to 20 barns, except for water
- Lighter nuclei slow neutrons down faster
- Low absorption is favorable – Hence why heavy water is useful

Solid Angle Review



$$d\Omega = \frac{dA}{r^2}$$

$$d\Omega = \sin \theta \, d\theta \, d\phi$$



“NE:150 Spring 2018 Lecture 19: Neutron Diffusion Equation – 3”, J. Vujic

Angular Flux

The path length per unit volume about $\vec{r}$ passed by neutrons with energies in dE about E at time t .

$$\psi(\vec{r}, E, \hat{\Omega}, t) \equiv vn(\vec{r}, E, \hat{\Omega}, t)$$

$$\text{units: } \frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV} \cdot \text{steradian}}$$

Scalar Flux

The number of neutrons penetrating a sphere with a cross sectional area of 1 cm^2 at $\vec{r}$, with energies in dE about E at time t .

$$\phi(\vec{r}, E, t) \equiv vn(\vec{r}, E, t)$$

$$\phi(\vec{r}, E, t) = \int_{4\pi} \psi(\vec{r}, E, \hat{\Omega}, t) d\hat{\Omega}$$

$$\phi(\vec{r}, E, t) = \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta \psi(\vec{r}, E, \hat{\Omega}, t)$$

$$\text{units: } \frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV}}$$

Net Current

Net number of particles crossing a unit area per second along a **direction normal** to that area with energies in $[E, E+dE]$ at time t .

$$\vec{J}(\vec{r}, E, t) = \int_{4\pi} \hat{\Omega} \psi(\vec{r}, E, \hat{\Omega}, t) d\hat{\Omega}$$
$$\vec{J}(\vec{r}, E, t) = \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta \cos\theta \psi(\vec{r}, E, \hat{\Omega}, t)$$

$$\text{units: } \frac{\text{neutrons}}{\text{cm}^2 \cdot \text{s} \cdot \text{MeV}}$$